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## Eigenvalues + Eigenvectors

linear operator  $A$

$|\phi\rangle$  is an eigenvector of  $A$  if

$$A|\phi\rangle = \lambda|\phi\rangle$$

scalar  $\lambda$  is called  
an eigenvalue of  $A$ .

If  $A$  is linear operator acting on vectors in  $\mathbb{C}^N$ ,  
then there is a set of eigenvectors that  
forms an orthonormal basis of  $\mathbb{C}^N$ .

$\{|\phi_1\rangle, |\phi_2\rangle, \dots, |\phi_N\rangle\}$  eigenvectors of  $A$

also an orthonormal basis of  $\mathbb{C}^N$ .

$$i \neq j \quad \langle \phi_i | A | \phi_j \rangle = \langle \phi_i | \lambda_j | \phi_j \rangle = \lambda_j \langle \phi_i | \phi_j \rangle$$

$$\langle \phi_i | A | \phi_i \rangle = \langle \phi_i | \lambda_i | \phi_i \rangle = \lambda_i \langle \phi_i | \phi_i \rangle = \lambda_i$$

When we express  $A$  using basis  $\{|\phi_i\rangle\}$

$$\sum_{j,k} \langle \phi_j | A | \phi_k \rangle |\phi_j\rangle \langle \phi_k| = \sum_{j=1}^N \lambda_j |\phi_j\rangle \langle \phi_j|$$

Theorem:  $A$  has a diagonal representation iff  $A$  is normal:  $AA^\dagger = A^\dagger A$ .

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Example:  $X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$

eigenval: 1                      -1

Eigenvektors  $|+\rangle$  and  $|-\rangle$

$$|+\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$|-\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

$$|+\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$|-\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} -1 \\ 1 \end{pmatrix}$$

$$X = \underbrace{|+\rangle\langle+|}_{\frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}} - \underbrace{|-\rangle\langle-|}_{\frac{1}{2} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}}$$

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If  $|v_i\rangle$  and  $|v_j\rangle$  are eigenvectors of  $A$  with different eigenvalues ( $\lambda_i \neq \lambda_j$ ), then  $\langle v_i | v_j \rangle = 0$ .

Non-degenerate case

Example:  $A$  is a  $4 \times 4$  matrix

$$\lambda_1 > \lambda_2 > \lambda_3 > \lambda_4$$

(all distinct)

$$\lambda_1 \quad |v_1\rangle$$

$$\lambda_2 \quad |v_2\rangle$$

$$\lambda_3 \quad |v_3\rangle$$

$$\lambda_4 \quad |v_4\rangle$$

If  $|v_i\rangle$ 's are normalized then the eigenvectors are all unique.

and they form an orthonormal basis

Degenerate Case.

$$\lambda_1 > \lambda_2 = \lambda_3 > \lambda_4$$

$$\lambda_1 - |v_1\rangle$$

$$\lambda_2 = \lambda_3 \begin{cases} |v_2\rangle \\ |v_3\rangle \end{cases}$$

$$\lambda_4 - |v_4\rangle$$

$$|v_1\rangle \quad |w_2\rangle \quad |w_3\rangle \quad |v_4\rangle$$

Are an orthonormal basis of eigenvectors if:

$|w_2\rangle$  and  $|w_3\rangle$  form an orthonormal basis of the space spanned by  $|v_2\rangle + |v_3\rangle$

$$\begin{matrix} |v_2\rangle & |v_3\rangle & \langle v_2 | v_3 \rangle = 0 \\ \frac{1}{\sqrt{2}} (|v_2\rangle + |v_3\rangle) & \frac{1}{\sqrt{2}} (|v_2\rangle - |v_3\rangle) \end{matrix}$$

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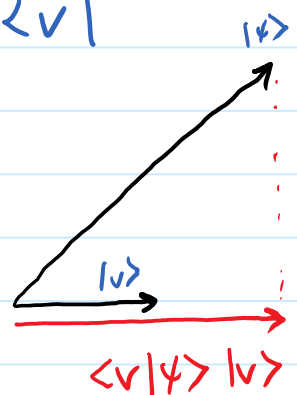
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Projectors: (Special class of linear operators in which eigenvalues = 1 or 0).

Vector  $|v\rangle$  (Normalized:  $\langle v|v\rangle = 1$ ).

Projector onto  $|v\rangle$ :  $P_v = |v\rangle\langle v|$

$$P_v|\psi\rangle = \langle v|\psi\rangle|v\rangle$$



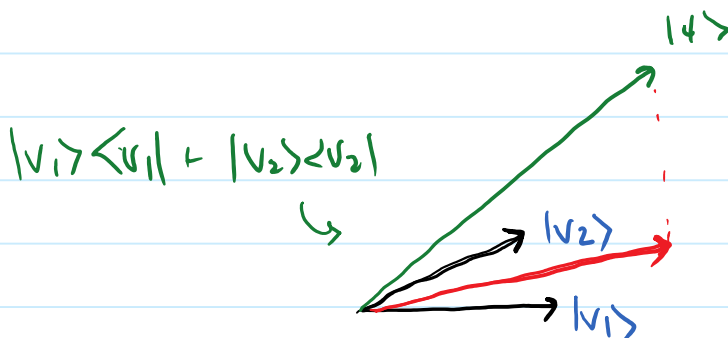
Can also project on to subspaces of higher dimension:

$$|v_1\rangle \dots |v_r\rangle \in \mathbb{C}^N$$

$$P = \sum_{j=1}^r |v_j\rangle\langle v_j|$$

$\langle v_j|v_k\rangle = \delta_{kj}$   
orthonormal set  
(generally does NOT span  $\mathbb{C}^N$ ).

$\delta_{kj} = 1$  if  $k=j$   
 $\delta_{kj} = 0$  if  $k \neq j$ .



Note:  
 $P^2 = P$ .

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## Measurement revisited

Measurement consists of an orthonormal basis  $\{|v_j\rangle\}$  and a real value  $r_j$  for each  $|v_j\rangle$ .

(Assume for now that  $r_j$ 's are all distinct)

Measure state  $|\phi\rangle$  in basis  $\{|v_j\rangle\}$

Outcome is  $r_j$  with prob  $|\langle v_j|\phi\rangle|^2$

$$|\phi\rangle = \alpha_1|v_1\rangle + \alpha_2|v_2\rangle + \dots + \alpha_n|v_n\rangle \quad \langle v_j|\phi\rangle = \alpha_j$$

$\langle v_j|$   $\alpha_1\langle v_j|v_1\rangle + \alpha_2\langle v_j|v_2\rangle \dots \alpha_j\langle v_j|v_j\rangle$

Afterwards state is  $|v_j\rangle$ .

qubit example:  $|\phi\rangle = \alpha|0\rangle + \beta|1\rangle$

Measure in  $|0\rangle^0 |1\rangle^1$  basis.

Probability outcome is 0 :  $|\langle 0|\phi\rangle|^2$

$$|\alpha|^2 + |\beta|^2 = 1. \quad = |\alpha|^2$$

Afterwards, state is  $|0\rangle$ .

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Could also measure  $|\phi\rangle = \alpha|0\rangle + \beta|1\rangle$

in the  $|+\rangle$   $|-\rangle$  basis.

Outcome is "+" with prob.  $|\langle +|\phi\rangle|^2$

$$\begin{aligned} |+\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) & \left( \underbrace{\langle +|\frac{1}{\sqrt{2}}}_{\langle +|} + \underbrace{\langle +|\frac{1}{\sqrt{2}}}_{\langle +|} \right) & \left( \underbrace{\alpha|0\rangle + \beta|1\rangle}_{|\phi\rangle} \right) \\ & & & \\ & & & = \frac{\beta}{\sqrt{2}} + \frac{\alpha}{\sqrt{2}} = \frac{\beta + \alpha}{\sqrt{2}} \end{aligned}$$

~~$\frac{\alpha}{\sqrt{2}} \langle +|0\rangle$~~

prob  $|+\rangle$   $|\langle +|\phi\rangle|^2 = \left| \frac{\beta + \alpha}{\sqrt{2}} \right|^2 \rightsquigarrow |+\rangle$

prob  $|-\rangle$   $|\langle -|\phi\rangle|^2 = \left| -\frac{\beta - \alpha}{\sqrt{2}} \right|^2 \rightsquigarrow |-\rangle$

State after measurement.

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3 qubit example:

$$\alpha_0 |000\rangle + \alpha_1 |001\rangle + \alpha_2 |010\rangle + \alpha_3 |011\rangle \\ + \alpha_4 |100\rangle + \alpha_5 |101\rangle + \alpha_6 |110\rangle + \alpha_7 |111\rangle$$

Measure all 3 qubits in the standard basis:

$ v_0\rangle =  000\rangle$	$r_0 = 0$
$ v_1\rangle =  001\rangle$	$r_1 = 1$
$ v_2\rangle =  010\rangle$	$r_2 = 2$
$ v_3\rangle =  011\rangle$	$r_3 = 3$
$ v_4\rangle =  100\rangle$	$r_4 = 4$
$ v_5\rangle =  101\rangle$	$r_5 = 5$
$ v_6\rangle =  110\rangle$	$r_6 = 6$
$ v_7\rangle =  111\rangle$	$r_7 = 7$

Outcome is  $j$   
w.p.  $|\langle v_j | \phi \rangle|^2 = |\alpha_j|^2$

Afterwards the state  
is  $|v_j\rangle$

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Measure just the first qubit:

degenerate case  
 $r_i$ 's not all distinct.

$$\begin{aligned} |V_0\rangle &= |000\rangle \\ |V_1\rangle &= |001\rangle \\ |V_2\rangle &= |010\rangle \\ |V_3\rangle &= |011\rangle \end{aligned} \quad r_0 = 0$$

$$P_0 = |V_0\rangle\langle V_0| + |V_1\rangle\langle V_1| + |V_2\rangle\langle V_2| + |V_3\rangle\langle V_3|$$

Projector onto the subspace of all states that have "0" as first qubit.

$$\begin{aligned} |V_4\rangle &= |100\rangle \\ |V_5\rangle &= |101\rangle \\ |V_6\rangle &= |110\rangle \\ |V_7\rangle &= |111\rangle \end{aligned} \quad r_1 = 1$$

$$P_1 = |V_4\rangle\langle V_4| + |V_5\rangle\langle V_5| + |V_6\rangle\langle V_6| + |V_7\rangle\langle V_7|$$

State:  $|\phi\rangle$  outcome is 0 w.p.  $|P_0|\phi\rangle|^2$

afterwards state is  $\frac{P_0|\phi\rangle}{|P_0|\phi\rangle}$

$$P_0|\phi\rangle =$$

$$(|V_0\rangle\langle V_0| + |V_1\rangle\langle V_1| + |V_2\rangle\langle V_2| + |V_3\rangle\langle V_3|)$$

$$\alpha_0|V_0\rangle + \alpha_1|V_1\rangle + \alpha_2|V_2\rangle + \alpha_3|V_3\rangle + \alpha_4|V_4\rangle + \alpha_5|V_5\rangle + \alpha_6|V_6\rangle + \alpha_7|V_7\rangle$$

$$\alpha_0 \langle V_0|V_0\rangle |V_0\rangle + \alpha_1|V_1\rangle + \alpha_2|V_2\rangle + \alpha_3|V_3\rangle$$

$$|P_0|\phi\rangle|^2 = |\alpha_0|^2 + |\alpha_1|^2 + |\alpha_2|^2 + |\alpha_3|^2$$

Final state is  $\frac{\alpha_0|V_0\rangle + \alpha_1|V_1\rangle + \alpha_2|V_2\rangle + \alpha_3|V_3\rangle}{(|\alpha_0|^2 + |\alpha_1|^2 + |\alpha_2|^2 + |\alpha_3|^2)^{1/2}}$



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## Generalized Measurement (degenerate case)

Orthogonal basis  $|v_1\rangle \dots |v_N\rangle$   
outcomes  $r_1 \dots r_N$  (real)

$r_j$ 's not all distinct.

$|v_1\rangle \dots |v_m\rangle$  all have  $r_j = R$   $1 \leq j \leq m$ .

$$P_R = \sum_{j=1}^m |v_j\rangle \langle v_j|$$

projector onto subspace spanned by  $|v_1\rangle \dots |v_m\rangle$

Outcome of measurement =  $R$  with prob

$$|P_R|\phi\rangle|^2$$

Afterwards New state is  $\frac{P_R|\phi\rangle}{|P_R|\phi\rangle}$

$$|\phi\rangle = \alpha_1|v_1\rangle + \alpha_2|v_2\rangle + \dots + \alpha_m|v_m\rangle + \alpha_{m+1}|v_{m+1}\rangle + \dots + \alpha_N|v_N\rangle$$

$$\text{Prob. outcome} = R: |\alpha_1|^2 + |\alpha_2|^2 + \dots + |\alpha_m|^2$$

$$\text{State becomes } \frac{\alpha_1|v_1\rangle + \alpha_2|v_2\rangle + \dots + \alpha_m|v_m\rangle}{(|\alpha_1|^2 + |\alpha_2|^2 + \dots + |\alpha_m|^2)^{1/2}}$$

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## Postulate of Quantum Mechanics : Measurement

Measurement operator (combines  $v_j$ 's +  $\langle v_j$ 's)

$$M = \sum_j v_j |v_j\rangle\langle v_j|$$

$|v_j\rangle$ 's are eigenvectors of  $M$ .  
 $v_j$ 's are eigenvalues.

$M$  is Hermitian (eigenvalues are real)  $M = M^\dagger$

$\Rightarrow$  Every physically observable quantity:  
 (Energy, momentum, location, ...)  
 is associated with a Hermitian operator

The outcome of measuring the quantity  
 is an eigenvalue of the operator

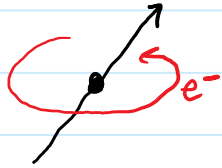
Afterwards the state  $|\phi\rangle$  becomes  $P|\phi\rangle$  (normalized)

$P$  is the projector onto the subspace  
 spanned by the eigenvectors whose  
 eigenvalues is the same as the outcome.

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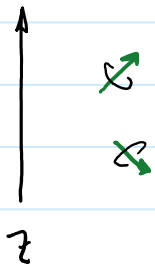
# Stern - Gerlach Experiment

1922 - Silver  
1927 - Hydrogen.



Atomic spin: Loop of "current" creates a magnetic dipole moment (bar magnet)

Experimental apparatus: direction of deflection depends on  $\hat{z}$  component of the atom's magnetic dipole moment.  $\hat{z}$  component is measured.

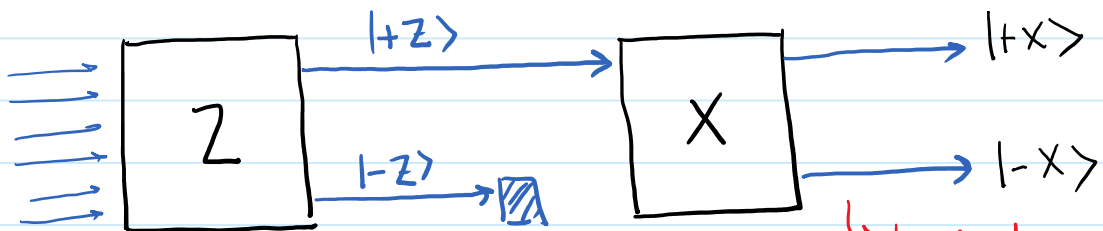


Would expect this to be uniformly distributed

Instead: produced two discrete beams corresponding to:  
↑ and ↓.

## Cascade two deflectors

Second turned on its side - measures spin in the  $\hat{x}$  direction.

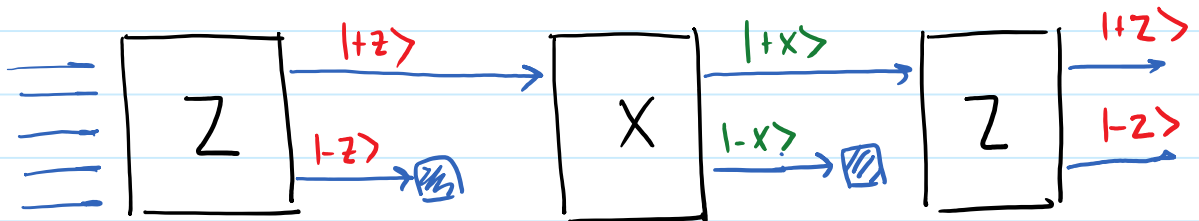


↳ Would have expected a single beam  $|+z\rangle$  has X-component = 0.

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Perhaps atoms have magnetic dipole moment along each axis independently



Quantum Explanation:  $|+z\rangle = |0\rangle$   $|-z\rangle = |1\rangle$

Z measures spin in  $|0\rangle$   $|1\rangle$  basis

X measures spin in  $|+\rangle$   $|-\rangle$  basis.